

## ADJOINT PROBLEM FOR THE LAPLACE EQUATION UNDER A NONLOCAL BOUNDARY CONDITION

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**ABSTRACT.** The paper is devoted to a new method of construction of the adjoint problem to a boundary value problem for an elliptic partial differential equation in a bounded domain with nonlocal boundary conditions. In this method, we do not use conditions except the given boundary ones. The method is universal in the sense that it can be applied both to linear ordinary differential equations and to linear partial differential equations.

**Keywords:** boundary value problem, adjoint problem, second Green's formula, sufficient boundary conditions.

**AMS Subject Classification:** 35J25.

### 1. INTRODUCTION

It is known that the Naimark method to construct the adjoint problem to a boundary value problem for a linear ordinary differential equation is based on the Lagrange formula related to the initial equation [8]. If the initial equation is of order  $n$ , one adds, to the given boundary conditions, some expressions with arbitrary coefficients so that one gets  $2n$  conditions. If these  $2n$  relations are linearly independent, they make it possible to determine the values of the unknown function as well as the values of its derivatives up to the order  $(n - 1)$  at the edges of the segment. The obtained values are substituted into the Lagrange formula and the result is assumed to be equal to zero. The boundary conditions of the adjoint problem include arbitrary coefficients which appeared in the added relations. Therefore, the boundary conditions of the adjoint problem are not unique. For linear ordinary differential equation of second order this difficulty was eliminated in [1] by a modification of the Naimark method. Later, this new method was applied to a linear ordinary differential equation of the third order [4]. The adjoint problem to a boundary value problem for an elliptic linear partial differential equation of the first order with nonlocal boundary condition has been constructed in a unique way (i.e. without ambiguous coefficients) [2]. The solution of the boundary value problem for an ordinary integro-differential equation with local boundary conditions was investigated in [3]. Note that the adjoint problem for elliptic partial differential equations of high order with local boundary conditions were considered in the works of Lions [5], [6], [7].

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## 2. FORMULATION OF THE PROBLEM

Let a bounded domain  $D$  in the  $(x_1, x_2)$ -plane be convex in the  $x_2$  direction,  $D = \{(x_1, x_2) : \gamma_1(x_1) \leq x_2 \leq \gamma_2(x_1)\}$ ,  $\Gamma = \partial D$ . This means, in particular, that it has a well-defined normal direction at each point.  $\gamma_1$  and  $\gamma_2$  are differentiable. Let us consider the following problem:

$$lu \equiv \Delta u(x) = \sum_{k=1}^2 \frac{\partial^2 u(x)}{\partial x_k^2} = 0, \quad x = (x_1, x_2) \in D \subset R^2, \quad (1)$$

$$\sum_{s=1}^2 \left[ \sum_{k=1}^2 \alpha_{mk}^{(s)}(x_1) \frac{\partial u(x)}{\partial x_k} + \alpha_m^{(s)}(x_1) u(x) \right] \Big|_{x_2=\gamma_s(x_1)} = 0, \quad m = 1, 2, \quad x_1 \in [a_1, b_1]. \quad (2)$$

Here,  $\alpha_{mk}^{(s)}(x_1)$  and  $\alpha_m^{(s)}(x_1)$ ,  $m, k, s = 1, 2, x_1 \in [a_1, b_1]$ , are continuous functions, conditions (2) are linearly independent. Let us construct the second Green formula for the equation (1)

$$(lu, v) = \sum_{k=1}^2 \int_D \frac{\partial^2 u(x)}{\partial x_k^2} \bar{v}(x) dx = \sum_{k=1}^2 \int_{\Gamma} \left[ \frac{\partial u(x)}{\partial x_k} \bar{v}(x) - u(x) \frac{\partial \bar{v}(x)}{\partial x_k} \right] \times \\ \times \cos(n, x_k) dx + (u, l^*v),$$

where  $n$ - is the external normal to  $\Gamma$ ,

$$l^*v = \Delta v(x) = \sum_{k=1}^2 \frac{\partial^2 v(x)}{\partial x_k^2}, \quad x \in D. \quad (3)$$

Let us transform the bilinear expression contained in Green's formula

$$B(u, v) = \sum_{k=1}^2 \int_{\Gamma} \left[ \frac{\partial u(x)}{\partial x_k} \bar{v}(x) - u(x) \frac{\partial \bar{v}(x)}{\partial x_k} \right] \cdot \cos(n, x_k) dx = \\ = \int_{a_1}^{b_1} \left\{ \left[ \frac{\partial u(x)}{\partial x_1} \bar{v}(x) - u(x) \frac{\partial \bar{v}(x)}{\partial x_1} \right] \Big|_{x_2=\gamma_1(x_1)} \cdot \cos(n, x_1) + \right. \\ \left. + \left[ \frac{\partial u(x)}{\partial x_2} \bar{v}(x) - u(x) \frac{\partial \bar{v}(x)}{\partial x_2} \right] \Big|_{x_2=\gamma_1(x_1)} \cdot \cos(n, x_2) + \right. \\ \left. + \left[ \frac{\partial u(x)}{\partial x_1} \bar{v}(x) - u(x) \frac{\partial \bar{v}(x)}{\partial x_1} \right] \Big|_{x_2=\gamma_2(x_1)} \cdot \cos(n, x_1) + \right. \\ \left. + \left[ \frac{\partial u(x)}{\partial x_2} \bar{v}(x) - u(x) \frac{\partial \bar{v}(x)}{\partial x_2} \right] \Big|_{x_2=\gamma_2(x_1)} \cdot \cos(n, x_2) \right\} \frac{dx_1}{\cos(x_1, \tau)} = 0. \quad (4)$$

Here  $\tau$  is the tangential direction to the boundary  $\Gamma$ .

One has:

$$u'(x_1, \gamma_k(x_1)) = \frac{\partial u(x)}{\partial x_1} \Big|_{x_2=\gamma_k(x_1)} + \frac{\partial u(x)}{\partial x_2} \Big|_{x_2=\gamma_k(x_1)} \gamma'_k(x_1), \quad k = 1, 2. \quad (5)$$

Thus, the expression (4) can be written as follows:

$$\begin{aligned}
 B(u, v) \equiv & \int_{a_1}^{b_1} \left\{ \bar{v}(x_1, \gamma_1(x_1)) \left[ u'(x_1, \gamma_1(x_1)) - \frac{\partial u(x)}{\partial x_2} \Big|_{x_2=\gamma_1(x_1)} \gamma'_{1}(x_1) \right] - \right. \\
 & \left. - u(x_1, \gamma_1(x_1)) \frac{\partial \bar{v}(x)}{\partial x_1} \Big|_{x_2=\gamma_1(x_1)} \right\} \gamma'_{1}(x_1) dx_1 - \\
 & - \int_{a_1}^{b_1} \left\{ \bar{v}(x_1, \gamma_1(x_1)) \frac{\partial u(x)}{\partial x_2} \Big|_{x_2=\gamma_1(x_1)} - u(x_1, \gamma_1(x_1)) \frac{\partial \bar{v}(x)}{\partial x_2} \Big|_{x_2=\gamma_1(x_1)} \right\} dx_1 - \\
 & - \int_{a_1}^{b_1} \left\{ \bar{v}(x_1, \gamma_2(x_1)) \left[ u'(x_1, \gamma_2(x_1)) - \frac{\partial u(x)}{\partial x_2} \Big|_{x_2=\gamma_2(x_1)} \gamma'_{2}(x_1) \right] - \right. \\
 & \left. - u(x_1, \gamma_2(x_1)) \frac{\partial \bar{v}(x)}{\partial x_1} \Big|_{x_2=\gamma_2(x_1)} \right\} \gamma'_{2}(x_1) dx_1 + \\
 & + \int_{a_1}^{b_1} \left\{ \bar{v}(x_1, \gamma_2(x_1)) \frac{\partial u(x)}{\partial x_2} \Big|_{x_2=\gamma_2(x_1)} - u(x_1, \gamma_2(x_1)) \frac{\partial \bar{v}(x)}{\partial x_2} \Big|_{x_2=\gamma_2(x_1)} \right\} dx_1 = \\
 & = \bar{v}(b_1, \gamma_1(b_1)) u(b_1, \gamma_1(b_1)) \cdot \gamma'_{1}(b_1) - \bar{v}(a_1, \gamma_1(a_1)) u(a_1, \gamma_1(a_1)) \cdot \gamma'_{1}(a_1) - \\
 & - \int_{a_1}^{b_1} u(x_1, \gamma_1(x_1)) [\gamma'_{1}(x_1) \bar{v}(x_1, \gamma_1(x_1))] ' dx_1 - \\
 & - \int_{a_1}^{b_1} \frac{\partial u(x)}{\partial x_2} \Big|_{x_2=\gamma_1(x_1)} \bar{v}(x_1, \gamma_1(x_1)) \gamma'^2_{1}(x_1) dx_1 - \int_{a_1}^{b_1} u(x_1, \gamma_1(x_1)) \frac{\partial \bar{v}(x)}{\partial x_1} \Big|_{x_2=\gamma_1(x_1)} \gamma'_{1}(x_1) dx_1 - \\
 & - \int_{a_1}^{b_1} \frac{\partial u(x)}{\partial x_2} \Big|_{x_2=\gamma_1(x_1)} \bar{v}(x_1, \gamma_1(x_1)) dx_1 + \int_{a_1}^{b_1} u(x_1, \gamma_1(x_1)) \frac{\partial \bar{v}(x)}{\partial x_2} \Big|_{x_2=\gamma_1(x_1)} dx_1 - \\
 & - \bar{v}(b_1, \gamma_2(b_1)) u(b_1, \gamma_2(b_1)) \cdot \gamma'_{2}(b_1) + \bar{v}(a_1, \gamma_2(a_1)) u(a_1, \gamma_2(a_1)) \cdot \gamma'_{2}(a_1) + \\
 & + \int_{a_1}^{b_1} u(x_1, \gamma_2(x_1)) [\gamma'_{2}(x_1) \bar{v}(x_1, \gamma_2(x_1))] ' dx_1 + \\
 & + \int_{a_1}^{b_1} \frac{\partial u(x)}{\partial x_2} \Big|_{x_2=\gamma_2(x_1)} \bar{v}(x_1, \gamma_2(x_1)) \gamma'^2_{2}(x_1) dx_1 + \int_{a_1}^{b_1} u(x_1, \gamma_2(x_1)) \frac{\partial \bar{v}(x)}{\partial x_1} \Big|_{x_2=\gamma_2(x_1)} \gamma'_{2}(x_1) dx_1 + \\
 & + \int_{a_1}^{b_1} \frac{\partial u(x)}{\partial x_2} \Big|_{x_2=\gamma_2(x_1)} \bar{v}(x_1, \gamma_2(x_1)) dx_1 - \int_{a_1}^{b_1} u(x_1, \gamma_2(x_1)) \frac{\partial \bar{v}(x)}{\partial x_2} \Big|_{x_2=\gamma_2(x_1)} dx_1 =
 \end{aligned}$$

$$\begin{aligned}
 &= \int_{a_1}^{b_1} \frac{\partial u(x)}{\partial x_2} \Big|_{x_2=\gamma_1(x_1)} [-\bar{v}(x_1, \gamma_1(x_1)) \gamma_1^{\prime 2}(x_1) - \bar{v}(x_1, \gamma_1(x_1))] dx_1 + \\
 &\quad + \int_{a_1}^{b_1} \frac{\partial u(x)}{\partial x_2} \Big|_{x_2=\gamma_2(x_1)} [\bar{v}(x_1, \gamma_2(x_1)) \gamma_2^{\prime 2}(x_1) + \bar{v}(x_1, \gamma_2(x_1))] dx_1 + \\
 &+ \int_{a_1}^{b_1} u(x_1, \gamma_1(x_1)) [-\bar{v}'(x_1, \gamma_1(x_1)) \gamma_1'(x_1) - \gamma_1''(x_1) \bar{v}(x_1, \gamma_1(x_1)) - \frac{\partial \bar{v}(x)}{\partial x_1} \Big|_{x_2=\gamma_1(x_1)} \gamma_1'(x_1) + \\
 &\quad + \frac{\partial \bar{v}(x)}{\partial x_2} \Big|_{x_2=\gamma_1(x_1)}] dx_1 + \int_{a_1}^{b_1} u(x_1, \gamma_2(x_1)) [-\bar{v}'(x_1, \gamma_2(x_1)) \gamma_2'(x_1) + \gamma_2''(x_1) \bar{v}(x_1, \gamma_2(x_1)) + \\
 &\quad + \frac{\partial \bar{v}(x)}{\partial x_1} \Big|_{x_2=\gamma_2(x_1)} \gamma_2'(x_1) - \frac{\partial \bar{v}(x)}{\partial x_2} \Big|_{x_2=\gamma_2(x_1)}] dx_1 = 0, \tag{6}
 \end{aligned}$$

here

$$u(b_1, \gamma_1(b_1)) = u(a_1, \gamma_1(a_1)) = u(a_1, \gamma_2(a_1)) = u(b_1, \gamma_2(b_1)) = 0. \tag{7}$$

Coming back to the boundary conditions (2), let us multiply them by the  $v_m(x_1)$ . We get

$$\begin{aligned}
 &\sum_{m=1}^2 \sum_{s=1}^2 \left[ \sum_{k=1}^2 \int_{a_1}^{b_1} \alpha_{mk}^{(s)}(x_1) \frac{\partial u(x)}{\partial x_k} \Big|_{x_2=\gamma_s(x_1)} \bar{v}_m(x_1) dx_1 + \right. \\
 &\quad \left. + \int_{a_1}^{b_1} \alpha_m^{(s)}(x_1) u(x_1, \gamma_s(x_1)) \bar{v}_m(x_1) dx_1 \right] = 0.
 \end{aligned}$$

Taking (5) into account, we get:

$$\begin{aligned}
 &\sum_{m=1}^2 \sum_{s=1}^2 \left\{ \int_{a_1}^{b_1} \alpha_{m1}^{(s)}(x_1) \bar{v}_m(x_1) \left[ u'(x_1, \gamma_s(x_1)) - \frac{\partial u(x)}{\partial x_2} \Big|_{x_2=\gamma_s(x_1)} \gamma_s'(x_1) \right] dx_1 + \right. \\
 &\quad \left. + \int_{a_1}^{b_1} \alpha_{m2}^{(s)}(x_1) \bar{v}_m(x_1) \frac{\partial u(x)}{\partial x_2} \Big|_{x_2=\gamma_s(x_1)} dx_1 + \int_{a_1}^{b_1} \alpha_m^{(s)}(x_1) u(x_1, \gamma_s(x_1)) \bar{v}_m(x_1) dx_1 \right\} = 0,
 \end{aligned}$$

or

$$\begin{aligned}
 &\sum_{s=1}^2 \sum_{m=1}^2 \int_{a_1}^{b_1} \frac{\partial u(x)}{\partial x_2} \Big|_{x_2=\gamma_s(x_1)} [\alpha_{m2}^{(s)}(x_1) \bar{v}_m(x_1) - \alpha_{m1}^{(s)}(x_1) \bar{v}_m(x_1) \gamma_s'(x_1)] dx_1 + \\
 &\quad + \sum_{s=1}^2 \sum_{m=1}^2 \int_{a_1}^{b_1} u(x_1, \gamma_s(x_1)) [\alpha_m^{(s)}(x_1) \bar{v}_m(x_1) - (\alpha_{m1}^{(s)}(x_1) \bar{v}_m(x_1))'] dx_1 = 0. \tag{8}
 \end{aligned}$$

Here the conditions (7) are applied.

**Remark 2.1.** The functions  $v_1(x_1)$  and  $v_2(x_1)$  can be chosen so that the coefficients in the

square brackets at two fixed functions from  $\frac{\partial u(x)}{\partial x_2} \Big|_{x_2=\gamma_1(x_1)}$ ,  $\frac{\partial u(x)}{\partial x_2} \Big|_{x_2=\gamma_s(x_1)}$ ,  $u(x_1, \gamma_1(x_1))$  and  $u(x_1, \gamma_2(x_1))$  coincide with the corresponding coefficients in (6). The integrals corresponding to these two summands are determined from (8) and substituted in (6). We obtain one expression depending on the remaining two functions from  $\frac{\partial u(x)}{\partial x_2} \Big|_{x_2=\gamma_1(x_1)}$ ,  $\frac{\partial u(x)}{\partial x_2} \Big|_{x_2=\gamma_s(x_1)}$ ,  $u(x_1, \gamma_1(x_1))$  and  $u(x_1, \gamma_2(x_1))$ . This expression is equal to zero independently of these two functions. Here we shall compare the coefficients at  $\frac{\partial u(x)}{\partial x_2} \Big|_{x_2=\gamma_1(x_1)}$  and  $u(x_1, \gamma_2(x_1))$ .

Thus, the functions  $v_m(x_1)$  are determined from the equations:

$$\sum_{m=1}^2 \left[ \alpha_{m2}^{(1)}(x_1) - \alpha_{m1}^{(1)}(x_1) \gamma'_{1}(x_1) \right] \bar{v}_m(x_1) = -\bar{v}(x_1, \gamma_1(x_1))\gamma'^2_{1}(x_1) - \bar{v}(x_1, \gamma_1(x_1)),$$

$$\sum_{m=1}^2 \left[ \alpha_m^{(2)}(x_1) \bar{v}_m(x_1) - \left( \alpha_{m1}^{(2)}(x_1) \bar{v}_m(x_1) \right)' \right] = \bar{v}'(x_1, \gamma_2(x_1))\gamma'_{2}(x_1) + \bar{v}(x_1, \gamma_2(x_1))\gamma''_{2}(x_1) +$$

$$+ \frac{\partial \bar{v}(x)}{\partial x_1} \Big|_{x_2=\gamma_2(x_1)} \gamma'_{2}(x_1) - \frac{\partial \bar{v}(x)}{\partial x_2} \Big|_{x_2=\gamma_2(x_1)}. \tag{9}$$

From the first equation under the condition

$$\alpha_{22}^{(1)}(x_1) - \alpha_{21}^{(1)}(x_1) \gamma'_{1}(x_1) \neq 0, \tag{10}$$

we obtain  $\bar{v}_1(x_1)$  and  $\bar{v}_2(x_1)$ :

$$\bar{v}_2(x_1) = -\frac{\alpha_{12}^{(1)}(x_1) - \alpha_{11}^{(1)}(x_1) \gamma'_{1}(x_1)}{\alpha_{22}^{(1)}(x_1) - \alpha_{21}^{(1)}(x_1) \gamma'_{1}(x_1)} \bar{v}_1(x_1) - \frac{[1 + \gamma'^2_{1}(x_1)] \bar{v}(x_1, \gamma_1(x_1))}{\alpha_{22}^{(1)}(x_1) - \alpha_{21}^{(1)}(x_1) \gamma'_{1}(x_1)}. \tag{11}$$

Substituting (11) in the second equation (9), we get:

$$\alpha_1^{(2)}(x_1) \bar{v}_1(x_1) - \alpha_2^{(2)}(x_1) \frac{[\alpha_{12}^{(1)}(x_1) - \alpha_{11}^{(1)}(x_1) \gamma'_{1}(x_1)] \bar{v}_1(x_1) + [1 + \gamma'^2_{1}(x_1)] \bar{v}(x_1, \gamma_1(x_1))}{\alpha_{22}^{(1)}(x_1) - \alpha_{21}^{(1)}(x_1) \gamma'_{1}(x_1)} -$$

$$- \left( \alpha_{11}^{(2)}(x_1) \bar{v}_1(x_1) \right)' +$$

$$+ \left\{ \alpha_{21}^{(2)}(x_1) \frac{[\alpha_{12}^{(1)}(x_1) - \alpha_{11}^{(1)}(x_1) \gamma'_{1}(x_1)] \bar{v}_1(x_1) + [1 + \gamma'^2_{1}(x_1)] \bar{v}(x_1, \gamma_1(x_1))}{\alpha_{22}^{(1)}(x_1) - \alpha_{21}^{(1)}(x_1) \gamma'_{1}(x_1)} \right\}' =$$

$$= \bar{v}'(x_1, \gamma_2(x_1)) \gamma'_{2}(x_1) +$$

$$+ \bar{v}(x_1, \gamma_2(x_1)) \gamma''_{2}(x_1) + \frac{\partial \bar{v}(x)}{\partial x_1} \Big|_{x_2=\gamma_2(x_1)} \gamma'_{2}(x_1) - \frac{\partial \bar{v}(x)}{\partial x_2} \Big|_{x_2=\gamma_2(x_1)}. \tag{12}$$

Thus, we obtain the linear equation (12) containing the first order derivative of the function  $\bar{v}_1(x_1)$ . The solution of this equation may be constructed in an analytic form. For this we need to know one solution of the non-homogeneous equation. The equation (12) may be written in the following way:

$$A(x_1) \bar{v}'_1(x_1) + B(x_1) \bar{v}_1(x_1) = F(x_1), \tag{13}$$

where

$$\begin{aligned}
A(x_1) &= -\alpha_{11}^{(2)}(x_1) + \alpha_{21}^{(2)}(x_1) \frac{\alpha_{12}^{(1)}(x_1) - \alpha_{11}^{(1)}(x_1) \gamma'_1(x_1)}{\alpha_{22}^{(1)}(x_1) - \alpha_{21}^{(1)}(x_1) \gamma'_1(x_1)}, \\
B(x_1) &= \alpha_1^{(2)}(x_1) - \alpha_2^{(2)}(x_1) \frac{\alpha_{12}^{(1)}(x_1) - \alpha_{11}^{(1)}(x_1) \gamma'_1(x_1)}{\alpha_{22}^{(1)}(x_1) - \alpha_{21}^{(1)}(x_1) \gamma'_1(x_1)} - \alpha_{11}^{(2)}(x_1) + \\
&\quad + \left\{ \alpha_{21}^{(2)}(x_1) \frac{\alpha_{12}^{(1)}(x_1) - \alpha_{11}^{(1)}(x_1) \gamma'_1(x_1)}{\alpha_{22}^{(1)}(x_1) - \alpha_{21}^{(1)}(x_1) \gamma'_1(x_1)} \right\}', \\
F(x_1) &= \alpha_2^{(2)}(x_1) \frac{[1 + \gamma_1'^2(x_1)] \bar{v}(x_1, \gamma_1(x_1))}{\alpha_{22}^{(1)}(x_1) - \alpha_{21}^{(1)}(x_1) \gamma'_1(x_1)} - \left\{ \alpha_{21}^{(2)}(x_1) \frac{[1 + \gamma_1'^2(x_1)] \bar{v}(x_1, \gamma_1(x_1))}{\alpha_{22}^{(1)}(x_1) - \alpha_{21}^{(1)}(x_1) \gamma'_1(x_1)} \right\}' - \\
&\quad \bar{v}'(x_1, \gamma_2(x_1)) \gamma'_2(x_1) + \bar{v}(x_1, \gamma_2(x_1)) \gamma''_2(x_1) + \left. \frac{\partial \bar{v}(x)}{\partial x_1} \right|_{x_2=\gamma_2(x_1)} \gamma'_2(x_1) - \left. \frac{\partial \bar{v}(x)}{\partial x_2} \right|_{x_2=\gamma_2(x_1)}.
\end{aligned}$$

If we have

$$A(x_1) \neq 0, \quad (14)$$

then we get (13):

$$\bar{v}'_1(x_1) = -\frac{B(x_1)}{A(x_1)} \bar{v}_1(x_1) + \frac{F(x_1)}{A(x_1)},$$

where

$$\bar{v}_1(x_1) = c e^{-\int_{a_1}^{x_1} \frac{B(t)}{A(t)} dt} + \int_{a_1}^{x_1} \frac{F(\tau)}{A(\tau)} e^{-\int_{\tau}^{x_1} \frac{B(t)}{A(t)} dt} d\tau,$$

is the general solution of (13). Assuming  $c = 0$  one gets the solution

$$\bar{v}_1(x_1) = \int_{a_1}^{x_1} \frac{F(\tau)}{A(\tau)} e^{-\int_{\tau}^{x_1} \frac{B(t)}{A(t)} dt} d\tau. \quad (15)$$

If we substitute this expression in (11), we get:

$$\begin{aligned}
\bar{v}_2(x_1) &= -\frac{\alpha_{12}^{(1)}(x_1) - \alpha_{11}^{(1)}(x_1) \gamma'_1(x_1)}{\alpha_{22}^{(1)}(x_1) - \alpha_{21}^{(1)}(x_1) \gamma'_1(x_1)} \int_{a_1}^{x_1} \frac{F(\tau)}{A(\tau)} e^{-\int_{\tau}^{x_1} \frac{B(t)}{A(t)} dt} d\tau - \\
&\quad - \frac{[1 + \gamma_1'^2(x_1)] \bar{v}(x_1, \gamma_1(x_1))}{\alpha_{22}^{(1)}(x_1) - \alpha_{21}^{(1)}(x_1) \gamma'_1(x_1)}.
\end{aligned} \quad (16)$$

Taking into account (9) and substituting (15) and (16) in (8), we get :

$$\begin{aligned}
& - \int_{a_1}^{b_1} (1 + \gamma_1'^2(x_1)) \bar{v}(x_1, \gamma_1(x_1)) \left. \frac{\partial u(x)}{\partial x_2} \right|_{x_2=\gamma_1(x_1)} dx_1 + \\
& + \int_{a_1}^{b_1} [\bar{v}'(x_1, \gamma_2(x_1)) \gamma'_2(x_1) + \bar{v}(x_1, \gamma_2(x_1)) \gamma''_2(x_1) +
\end{aligned}$$

$$\begin{aligned}
 & + \left. \frac{\partial \bar{v}(x)}{\partial x_1} \right|_{x_2=\gamma_2(x_1)} \gamma'_{t_2}(x_1) - \left. \frac{\partial \bar{v}(x)}{\partial x_2} \right|_{x_2=\gamma_2(x_1)} \Big] u(x_1, \gamma_2(x_1)) dx_1 = \\
 & = - \sum_{m=1}^2 \int_{a_1}^{b_1} \left[ \alpha_{m2}^{(2)}(x_1) - \alpha_{m1}^{(2)}(x_1) \gamma'_{t_2}(x_1) \right] \bar{v}_m(x_1) \left. \frac{\partial u(x)}{\partial x_2} \right|_{x_2=\gamma_2(x_1)} dx_1 - \\
 & - \sum_{m=1}^2 \int_{a_1}^{b_1} \left[ \alpha_m^{(1)}(x_1) \bar{v}_m(x_1) - \left( \alpha_{m1}^{(1)}(x_1) \bar{v}_m(x_1) \right)' \right] u(x_1, \gamma_1(x_1)) dx_1.
 \end{aligned}$$

Finally, substituting this expression in (6), we get:

$$\begin{aligned}
 B(u, v) = & - \sum_{m=1}^2 \int_{a_1}^{b_1} \left[ \alpha_{m2}^{(2)}(x_1) - \alpha_{m1}^{(2)}(x_1) \gamma'_{t_2}(x_1) \right] \bar{v}_m(x_1) \left. \frac{\partial u(x)}{\partial x_2} \right|_{x_2=\gamma_2(x_1)} dx_1 - \\
 & - \sum_{m=1}^2 \int_{a_1}^{b_1} \left[ \alpha_m^{(1)}(x_1) \bar{v}_m(x_1) - \left( \alpha_{m1}^{(1)}(x_1) \bar{v}_m(x_1) \right)' \right] u(x_1, \gamma_1(x_1)) dx_1 + \\
 & + \int_{a_1}^{b_1} (1 + \gamma_{t_2}^2(x_1)) \bar{v}(x_1, \gamma_2(x_1)) \left. \frac{\partial u(x)}{\partial x_2} \right|_{x_2=\gamma_2(x_1)} dx_1 + \\
 & + \int_{a_1}^{b_1} \left[ -\bar{v}'(x_1, \gamma_1(x_1)) \gamma'_{t_1}(x_1) - \bar{v}(x_1, \gamma_1(x_1)) \gamma''_{t_1}(x_1) - \right. \\
 & \left. - \left. \frac{\partial \bar{v}(x)}{\partial x_1} \right|_{x_2=\gamma_1(x_1)} \gamma'_{t_1}(x_1) + \left. \frac{\partial \bar{v}(x)}{\partial x_2} \right|_{x_2=\gamma_1(x_1)} \right] u(x_1, \gamma_1(x_1)) dx_1 = 0.
 \end{aligned}$$

Since the equation holds for arbitrary  $\left. \frac{\partial u(x)}{\partial x_2} \right|_{x_2=\gamma_2(x_1)}$  and  $u(x_1, \gamma_1(x_1))$ , we get the following boundary condition for the adjoint problem:

$$\begin{aligned}
 & - \sum_{m=1}^2 \left[ \bar{\alpha}_{m2}^{(2)}(x_1) - \bar{\alpha}_{m1}^{(2)}(x_1) \gamma'_{t_2}(x_1) \right] v_m(x_1) + (1 + \gamma_{t_2}^2(x_1)) v(x_1, \gamma_2(x_1)) = 0, \\
 & - \sum_{m=1}^2 \left[ \bar{\alpha}_m^{(1)}(x_1) v_m(x_1) - \left( \bar{\alpha}_{m1}^{(1)}(x_1) v_m(x_1) \right)' \right] - v'(x_1, \gamma_1(x_1)) \gamma'_{t_1}(x_1) - \\
 & - v(x_1, \gamma_1(x_1)) \gamma''_{t_1}(x_1) - \left. \frac{\partial v(x)}{\partial x_1} \right|_{x_2=\gamma_1(x_1)} \gamma'_{t_1}(x_1) + \left. \frac{\partial v(x)}{\partial x_2} \right|_{x_2=\gamma_1(x_1)} = 0. \tag{17}
 \end{aligned}$$

Substituting  $v_1(x_1)$  and  $v_2(x_1)$  from (15) and (16) in (17), we get:

$$\begin{aligned}
 & - \left[ \bar{\alpha}_{12}^{(2)}(x_1) - \bar{\alpha}_{11}^{(2)}(x_1) \gamma'_{t_2}(x_1) \right] \int_{a_1}^{x_1} \frac{\bar{F}(\tau)}{\bar{A}(\tau)} e^{-\int_{\tau}^{x_1} \frac{\bar{B}(t)}{\bar{A}(t)} dt} d\tau + \\
 & + \left[ \bar{\alpha}_{22}^{(2)}(x_1) - \bar{\alpha}_{21}^{(2)}(x_1) \gamma'_{t_2}(x_1) \right] \left\{ \frac{[1 + \gamma_{t_1}^2(x_1)] v(x_1, \gamma_1(x_1))}{\bar{\alpha}_{22}^{(1)}(x_1) - \bar{\alpha}_{21}^{(1)}(x_1) \gamma'_{t_1}(x_1)} + \right. \\
 & \left. + \frac{\bar{\alpha}_{12}^{(1)}(x_1) - \bar{\alpha}_{11}^{(1)}(x_1) \gamma'_{t_1}(x_1)}{\bar{\alpha}_{22}^{(1)}(x_1) - \bar{\alpha}_{21}^{(1)}(x_1) \gamma'_{t_1}(x_1)} \int_{a_1}^{x_1} \frac{\bar{F}(\tau)}{\bar{A}(\tau)} e^{-\int_{\tau}^{x_1} \frac{\bar{B}(t)}{\bar{A}(t)} dt} d\tau \right\} +
 \end{aligned}$$

$$\begin{aligned}
 & + [1 + \gamma'^2_2(x_1)] v(x_1, \gamma_2(x_1)) = 0, \tag{18} \\
 & -\bar{\alpha}^{(1)}_1(x_1) \int_{a_1}^{x_1} \frac{\bar{F}(\tau)}{\bar{A}(\tau)} e^{-\int_{\tau}^{x_1} \frac{\bar{B}(t)}{\bar{A}(t)} dt} d\tau + \bar{\alpha}^{(1)}_2(x_1) \left\{ \frac{[1 + \gamma'^2_1(x_1)] v(x_1, \gamma_1(x_1))}{\bar{\alpha}^{(1)}_{22}(x_1) - \bar{\alpha}^{(1)}_{21}(x_1) \gamma'_1(x_1)} + \right. \\
 & \quad \left. + \frac{\bar{\alpha}^{(1)}_{12}(x_1) - \bar{\alpha}^{(1)}_{11}(x_1) \gamma'_1(x_1)}{\bar{\alpha}^{(1)}_{22}(x_1) - \bar{\alpha}^{(1)}_{21}(x_1) \gamma'_1(x_1)} \int_{a_1}^{x_1} \frac{\bar{F}(\tau)}{\bar{A}(\tau)} e^{-\int_{\tau}^{x_1} \frac{\bar{B}(t)}{\bar{A}(t)} dt} d\tau \right\} + \\
 & + \left\{ \bar{\alpha}^{(1)}_{11}(x_1) \int_{a_1}^{x_1} \frac{\bar{F}(\tau)}{\bar{A}(\tau)} e^{-\int_{\tau}^{x_1} \frac{\bar{B}(t)}{\bar{A}(t)} dt} d\tau - \bar{\alpha}^{(1)}_{21}(x_1) \left[ \frac{[1 + \gamma'^2_1(x_1)] v(x_1, \gamma_1(x_1))}{\bar{\alpha}^{(1)}_{22}(x_1) - \bar{\alpha}^{(1)}_{21}(x_1) \gamma'_1(x_1)} + \right. \right. \\
 & \quad \left. \left. + \frac{\bar{\alpha}^{(1)}_{12}(x_1) - \bar{\alpha}^{(1)}_{11}(x_1) \gamma'_1(x_1)}{\bar{\alpha}^{(1)}_{22}(x_1) - \bar{\alpha}^{(1)}_{21}(x_1) \gamma'_1(x_1)} \int_{a_1}^{x_1} \frac{\bar{F}(\tau)}{\bar{A}(\tau)} e^{-\int_{\tau}^{x_1} \frac{\bar{B}(t)}{\bar{A}(t)} dt} d\tau \right] \right\}' - v'(x_1, \gamma_1(x_1)) \gamma'_1(x_1) - \\
 & - v(x_1, \gamma_1(x_1)) \gamma''_1(x_1) - \left. \frac{\partial v(x)}{\partial x_1} \right|_{x_2=\gamma_1(x_1)} \gamma'_1(x_1) + \left. \frac{\partial v(x)}{\partial x_2} \right|_{x_2=\gamma_1(x_1)} = 0. \tag{19}
 \end{aligned}$$

The conditions (18) and (19) can be rewritten in the following way:

$$\begin{aligned}
 & \frac{\bar{\alpha}^{(2)}_{22}(x_1) - \bar{\alpha}^{(2)}_{21}(x_1) \gamma'_2(x_1)}{\bar{\alpha}^{(1)}_{22}(x_1) - \bar{\alpha}^{(1)}_{21}(x_1) \gamma'_1(x_1)} [1 + \gamma'^2_1(x_1)] v(x_1, \gamma_1(x_1)) + [1 + \gamma'^2_2(x_1)] v(x_1, \gamma_2(x_1)) + \\
 & + \left\{ \frac{\bar{\alpha}^{(2)}_{22}(x_1) - \bar{\alpha}^{(2)}_{21}(x_1) \gamma'_2(x_1)}{\bar{\alpha}^{(1)}_{22}(x_1) - \bar{\alpha}^{(1)}_{21}(x_1) \gamma'_1(x_1)} [\bar{\alpha}^{(1)}_{12}(x_1) - \bar{\alpha}^{(1)}_{11}(x_1) \gamma'_1(x_1)] - \right. \\
 & \quad \left. - [\bar{\alpha}^{(2)}_{12}(x_1) - \bar{\alpha}^{(2)}_{11}(x_1) \gamma'_2(x_1)] \right\} \cdot \int_{a_1}^{x_1} e^{-\int_{\tau}^{x_1} \frac{\bar{B}(t)}{\bar{A}(t)} dt} \frac{d\tau}{\bar{A}(\tau)} \times \\
 & \times \left\{ \left[ \frac{\bar{\alpha}^{(2)}_2(\tau)}{\bar{\alpha}^{(1)}_{22}(\tau) - \bar{\alpha}^{(1)}_{21}(\tau) \gamma'_1(\tau)} [1 + \gamma'^2_1(\tau)] - \left( \frac{\bar{\alpha}^{(2)}_{21}(\tau)}{\bar{\alpha}^{(1)}_{22}(\tau) - \bar{\alpha}^{(1)}_{21}(\tau) \gamma'_1(\tau)} [1 + \gamma'^2_1(\tau)] \right) \right]' \times \right. \\
 & \quad \times v(\tau, \gamma_1(\tau)) + \gamma'_2(\tau) v'(\tau, \gamma_2(\tau)) - \left( \frac{\bar{\alpha}^{(2)}_{21}(\tau)}{\bar{\alpha}^{(1)}_{22}(\tau) - \bar{\alpha}^{(1)}_{21}(\tau) \gamma'_1(\tau)} [1 + \gamma'^2_1(\tau)] \right) \times \\
 & \quad \times v'(\tau, \gamma_1(\tau)) + \gamma''_2(\tau) v(\tau, \gamma_2(\tau)) + \left. \frac{\partial v(\tau, x_2)}{\partial \tau} \right|_{x_2=\gamma_2(\tau)} \gamma'_2(\tau) - \\
 & \quad \left. - \frac{\partial v(\tau, x_2)}{\partial x_2} \right|_{x_2=\gamma_2(\tau)} \Big\} = 0, \tag{20}
 \end{aligned}$$

$$\begin{aligned}
 & \left\{ \frac{\bar{\alpha}^{(1)}_2(x_1)}{\bar{\alpha}^{(1)}_{22}(x_1) - \bar{\alpha}^{(1)}_{21}(x_1) \gamma'_1(x_1)} [1 + \gamma'^2_1(x_1)] - \left( \frac{\bar{\alpha}^{(1)}_{21}(x_1)}{\bar{\alpha}^{(1)}_{22}(x_1) - \bar{\alpha}^{(1)}_{21}(x_1) \gamma'_1(x_1)} [1 + \gamma'^2_1(x_1)] \right) \right\}' - \\
 & - \gamma''_1(x_1) \Big\} v(x_1, \gamma_1(x_1)) - \left\{ \gamma'_1(x_1) + \frac{\bar{\alpha}^{(1)}_{21}(x_1)}{\bar{\alpha}^{(1)}_{22}(x_1) - \bar{\alpha}^{(1)}_{21}(x_1) \gamma'_1(x_1)} [1 + \gamma'^2_1(x_1)] \right\} \times
 \end{aligned}$$



$$\begin{aligned}
 & \times v'(x_1, \gamma_1(x_1)) + \left[ \bar{\alpha}_2^{(1)}(x_1) \frac{\bar{\alpha}_{12}^{(1)}(x_1) - \bar{\alpha}_{11}^{(1)}(x_1) \gamma'_{1}(x_1)}{\bar{\alpha}_{22}^{(1)}(x_1) - \bar{\alpha}_{21}^{(1)}(x_1) \gamma'_{1}(x_1)} - \bar{\alpha}_1^{(1)}(x_1) \right] \int_{a_1}^{x_1} e^{-\int_{\tau}^{x_1} \frac{\bar{B}(t)}{A(t)} dt} \frac{d\tau}{\bar{A}(\tau)} \times \\
 & \times \left\{ \left[ \frac{\bar{\alpha}_2^{(2)}(\tau)}{\bar{\alpha}_{22}^{(1)}(\tau) - \bar{\alpha}_{21}^{(1)}(\tau) \gamma'_{1}(\tau)} [1 + \gamma'^2_{1}(\tau)] - \left( \frac{\bar{\alpha}_{21}^{(2)}(\tau)}{\bar{\alpha}_{22}^{(1)}(\tau) - \bar{\alpha}_{21}^{(1)}(\tau) \gamma'_{1}(\tau)} [1 + \gamma'^2_{1}(\tau)] \right)' \right] \times \right. \\
 & \quad \times v(\tau, \gamma_1(\tau)) + \gamma'_{2}(\tau) v'(\tau, \gamma_2(\tau)) - \left( \frac{\bar{\alpha}_{21}^{(2)}(\tau)}{\bar{\alpha}_{22}^{(1)}(\tau) - \bar{\alpha}_{21}^{(1)}(\tau) \gamma'_{1}(\tau)} [1 + \gamma'^2_{1}(\tau)] \right) \times \\
 & \quad \times v'(\tau, \gamma_1(\tau)) + \gamma''_{2}(\tau) v(\tau, \gamma_2(\tau)) + \left. \frac{\partial v(\tau, x_2)}{\partial \tau} \Big|_{x_2=\gamma_2(\tau)} \gamma'_{2}(\tau) - \right. \\
 & \quad \left. - \frac{\partial v(\tau, x_2)}{\partial x_2} \Big|_{x_2=\gamma_2(\tau)} \right\} + \left\{ \left[ \bar{\alpha}_{11}^{(1)}(x_1) - \bar{\alpha}_{21}^{(1)}(x_1) \frac{\bar{\alpha}_{12}^{(1)}(x_1) - \bar{\alpha}_{11}^{(1)}(x_1) \gamma'_{1}(x_1)}{\bar{\alpha}_{22}^{(1)}(x_1) - \bar{\alpha}_{21}^{(1)}(x_1) \gamma'_{1}(x_1)} \right] \times \right. \\
 & \quad \times \int_{a_1}^{x_1} e^{-\int_{\tau}^{x_1} \frac{\bar{B}(t)}{A(t)} dt} \frac{d\tau}{\bar{A}(\tau)} \left\{ \left[ \frac{\bar{\alpha}_2^{(2)}(\tau)}{\bar{\alpha}_{22}^{(1)}(\tau) - \bar{\alpha}_{21}^{(1)}(\tau) \gamma'_{1}(\tau)} [1 + \gamma'^2_{1}(\tau)] - \right. \right. \\
 & \quad \left. \left. - \left( \frac{\bar{\alpha}_{21}^{(2)}(\tau)}{\bar{\alpha}_{22}^{(1)}(\tau) - \bar{\alpha}_{21}^{(1)}(\tau) \gamma'_{1}(\tau)} [1 + \gamma'^2_{1}(\tau)] \right)' \right] v(\tau, \gamma_1(\tau)) + \gamma'_{2}(\tau) v'(x_1, \gamma_2(\tau)) - \right. \\
 & \quad \left. - \left( \frac{\bar{\alpha}_{21}^{(2)}(\tau)}{\bar{\alpha}_{22}^{(1)}(\tau) - \bar{\alpha}_{21}^{(1)}(\tau) \gamma'_{1}(\tau)} [1 + \gamma'^2_{1}(\tau)] \right) \times v'(\tau, \gamma_1(\tau)) + \gamma''_{2}(\tau) v(\tau, \gamma_2(\tau)) + \right. \\
 & \quad \left. + \frac{\partial v(\tau, x_2)}{\partial \tau} \Big|_{x_2=\gamma_2(\tau)} \gamma'_{2}(\tau) - \frac{\partial v(\tau, x_2)}{\partial x_2} \Big|_{x_2=\gamma_2(\tau)} \right\} \Big\} - \frac{\partial v(x)}{\partial x_1} \Big|_{x_2=\gamma_1(x_1)} \gamma'_{1}(x_1) + \\
 & \quad + \frac{\partial v(x)}{\partial x_2} \Big|_{x_2=\gamma_1(x_1)} = 0. \tag{21}
 \end{aligned}$$

Thus, we obtain the following statement:

**Theorem 2.1.** *If the domain  $D = \{(x_1, x_2) : \gamma_1(x_1) \leq x_2 \leq \gamma_2(x_1)\}$  is bounded, the functions  $\gamma_1(x_1)$  and  $\gamma_2(x_1)$  are two times differentiable, the boundary conditions (2) are linear independent, the coefficients in these conditions satisfy*

$$\alpha_{mk}^{(1)}(x_1) \in C^{(1)}(a_1, b_1), \alpha_{m1}^{(2)}(x_1) \in C^{(1)}(a_1, b_1), m, k = 1, 2; \alpha_{m2}^{(2)}(x_1) \in C(a_1, b_1),$$

$$\alpha_m^{(s)}(x_1) \in C^{(1)}(a_1, b_1), m, s = 1, 2;$$

and (7), (10) and (14) hold, then the problem adjoint to (1)-(2) is given by (3), (20), (21).

**Remark 2.2.** From (20) and (21) one can see that the boundary conditions of the adjoint problem contain both non-local and global terms. In other words, these conditions contain  $v(x_1, \gamma_1(x_1)), v(x_1, \gamma_2(x_1)), \frac{\partial v(x)}{\partial x_1} \Big|_{x_2=\gamma_1(x_1)}, \frac{\partial v(x)}{\partial x_1} \Big|_{x_2=\gamma_2(x_1)}, \frac{\partial v(x)}{\partial x_2} \Big|_{x_2=\gamma_1(x_1)}, \frac{\partial v(x)}{\partial x_2} \Big|_{x_2=\gamma_2(x_1)}$  both in the integrand (global terms) and outside of the integrals (non-local terms).

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